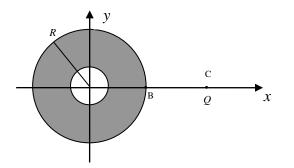
# FUNDAMENTALS OF ELECTROMAGNETIC FIELDS – January 22<sup>nd</sup>, 2018

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**Problem 1** - A uniform charge density  $\rho = 10^{-10}$  C/m<sup>3</sup> is embedded in a hollow dielectric sphere with radius R = 4 m (grey part in the figure). A point charge, with value  $Q = 3 \cdot 10^{-9}$  C, is fixed in point C(2*R*,0). Another charge *q*, free to move and with value -*Q*, is placed at the sphere center. Determine the value of another point charge *Q<sub>B</sub>* to be placed in B(*R*,0) such that *q* is in equilibrium.



### Solution

The electric field in (0,0) due to the charged hollow sphere is zero, as stated by the Gauss' theorem:

$$\int_{S} \vec{D} \cdot d\vec{S} = \int_{V} \rho \, dV$$

As q is negative and Q positive, the latter will tend to attract q; therefore, q will be subject to the following force:

$$\vec{F}_{C} = q\vec{E}_{Q}(0,0) = -|Q|\vec{E}_{Q}(0,0) = -|Q|\left[-\frac{|Q|}{4\pi\varepsilon_{0}(2R)^{2}}\vec{\mu}_{x}\right] = \frac{|Q|^{2}}{16\pi\varepsilon_{0}R^{2}}\vec{\mu}_{x}$$

In order to counterbalance the effect of Q in C, the charge to be placed in B must be negative so as to repulse q, which will be subject to:

$$\vec{F}_{B} = q\vec{E}_{Q_{B}}(0,0) = -|Q|\vec{E}_{Q_{B}}(0,0) = -|Q|\left[\frac{|Q_{B}|}{4\pi\varepsilon_{0}(R)^{2}}\vec{\mu}_{x}\right] = -\frac{|Q||Q_{B}|}{4\pi\varepsilon_{0}R^{2}}\vec{\mu}_{x}$$

The equilibrium for q is reached if:

 $\vec{F} = \vec{F}_B + \vec{F}_C = 0$ 

which yields:

$$\frac{|Q||Q_B|}{4\pi\varepsilon_0 R^2} = \frac{|Q|^2}{16\pi\varepsilon_0 R^2}$$

Therefore:

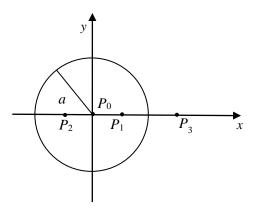
$$|Q_B| = \frac{|Q|}{4} = 7.5 \cdot 10^{-10} \text{ C} \Rightarrow Q_B = -7.5 \cdot 10^{-10} \text{ C}$$

**<u>Problem 2</u>** - Let us consider the following distribution for the current density flowing inside a cylindrical wire (see the wire section in the figure with radius a = 0.5 cm):

$$\vec{J} = \begin{cases} -\frac{2\rho}{a}\vec{a}_z & \rho \le a\\ 0 & \rho > a \end{cases} \quad A/m^2$$

where  $\rho$  is the distance from the wire axis.

Calculate the magnetic field (full vector) in the following points  $P_0$  (0,0,0),  $P_1(a/2,0,0)$ ,  $P_2(-a/2,0,0)$  e  $P_3(2a,0,0)$ .



## Solution

The solution is easily achieved using the Ampère's theorem, i.e.:

$$\int_{l} \vec{H} \cdot d\vec{l} = \int_{S} \vec{J} \cdot d\vec{S}$$

For point  $P_0$ ,  $\vec{H}(P_0) = 0$  because no current is included in the circle *l*.

For point  $P_1$ , Ampère's theorem becomes:

$$\left|\vec{H}\right| 2\pi\rho = \int_{0}^{\rho} \frac{2r}{a} 2\pi r \, dr = \frac{4\pi}{3a} \rho^{3} \implies \vec{H}(P_{1}) = -\frac{2\rho^{2}}{3a} \vec{a}_{y} \Big|_{P_{1}} = -8.3 \cdot 10^{-4} \vec{a}_{y} \text{ A/m}$$

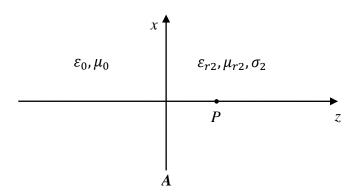
For point  $P_2$ , the solution is the same as for point P1, but the magnetic field vector changes direction:

$$\vec{H}(P_1) = 8.3 \cdot 10^{-4} \vec{a}_y$$
 A/m

Finally, for point *P*<sub>3</sub>:

$$\left|\vec{H}\right| 2\pi\rho = \int_{0}^{a} \frac{2r}{a} 2\pi r \, dr = \frac{4\pi}{3} a^{2} \implies \vec{H} = -\frac{2a^{2}}{3\rho} \vec{a}_{y} \Big|_{P_{3}} = -0.0017 \vec{a}_{y} \text{ A/m}$$

**Problem 3** - A uniform plane wave (frequency f = 200 MHz) propagates in free space and impinges on a dielectric material ( $\varepsilon_{r2} = 4$ ,  $\mu_{r2} = 1$ ,  $\sigma_2 = 4.5 \cdot 10^{-2}$  S/m). The electric field of the incident wave at (0,0,0) is  $\vec{E}_i(0,0,0) = -j\vec{a}_x$  V/m. Calculate the electric field in  $P(0,0,\lambda_2)$  and the power density carried by the wave in  $P(0,0,\lambda_2)$  ( $\lambda_2$  is the wavelength in the second medium).



#### Solution

First, let us calculate the intrinsic impedance for the two media. For the first one:

 $\eta_1 = \eta_0 \approx 377 \ \Omega$ 

For the second one, the loss tangent is  $\sigma/\omega\varepsilon \approx 1$ . Therefore no accurate approximations are possible:

$$\eta_2 = \sqrt{\frac{j\omega\mu_2}{\sigma + j\omega\varepsilon_2}} = 146 + j61\Omega$$

The reflection coefficient is:

$$\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = -0.423 + j0.166$$

Also, let us calculate the propagation constant for the second medium:

$$\gamma_2 = \alpha_2 + j\beta_2 = \sqrt{(j\omega\mu_2)(\sigma + j\omega\epsilon_2)} = 3.85 + j9.23$$
 1/m

The wavelength is obtained from  $\beta_2$  as:

$$\lambda_2 = \frac{2\pi}{\beta_2} \approx 0.68 \text{ m}$$

The field transmitted into the second medium is:

$$\vec{E}_2(0,0,0) = \vec{E}_i(0,0,0)(1+\Gamma) = (0.166 - j0.577)\vec{\mu}_x$$
 V/m

The electric field in *P* is therefore:

$$\vec{E}_2(P) = \vec{E}_2(0,0,0)e^{-\gamma_2\lambda_2} = (0.012 - j0.042)\vec{\mu}_x$$
 V/m

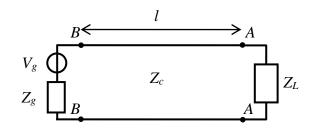
Finally, the power density is:

$$S(P) = \frac{1}{2} \frac{\left| \vec{E}_2(P) \right|^2}{\left| \eta_2 \right|} \cos(\angle \eta_2) = 5.5 \ \mu W/m^2$$

**Problem 4** - A source with voltage  $V_g = 50$  V and internal impedance  $Z_g = 100 \Omega$  is connected to a load  $Z_L = 150 \Omega$  by a transmission line with characteristic impedance  $Z_C = 50 \Omega$ . The line length is l = 4 m and the frequency is f = 300 MHz.

Calculate:

- a) The power absorbed by the load
- b) The voltage at the beginning of the line (section BB below),  $V_B$
- c) The trend of  $V_B$  in time



### Solution

a) The wavelength is:

 $\lambda = \lambda_0 = c/f = 1$  m The reflection coefficient at section AA is:

$$\Gamma_A = \frac{Z_L - Z_C}{Z_L + Z_C} = 0.5$$

The reflection coefficient at section BB is:

$$\Gamma_{B} = \Gamma_{A} \mathbf{e}^{-j2\beta l} = \Gamma_{A} \mathbf{e}^{-j2\left(\frac{2\pi}{\lambda}\right)4\lambda} = \Gamma_{A} \mathbf{e}^{-j16\pi} = 0.5$$

Therefore, the input impedance is:

$$Z_B = Z_C \frac{1 + \Gamma_B}{1 + \Gamma_B} = 150 \quad \Omega$$

The reflection coefficient for the source is:

$$\Gamma_g = \frac{Z_B - Z_g}{Z_B + Z_g} = 0.2$$

Therefore, the power crossing section BB, i.e. reaching the load is:

$$P_L = P_{AV} (1 - |\Gamma_g|^2) = 3 \text{ W}$$

## b) The voltage at the beginning of the line is:

$$V_B = V_g \frac{Z_B}{Z_B + Z_g} = 30 \text{ V}$$

c) The trend of  $V_B$  in time is given by:

$$v_B(t) = \operatorname{Re}\left[V_B e^{j\omega t}\right] = 30\cos(\omega t) \quad \nabla$$